

LEPの精密測定と新しい物理

2001. 11. 5

LEP ツンホツウム @ 東大

KEK 萩原 薫

• $N_\nu = 3$

• 弱い相互作用のゲージ理論を検証

• 大統一理論・超対称性のヒント

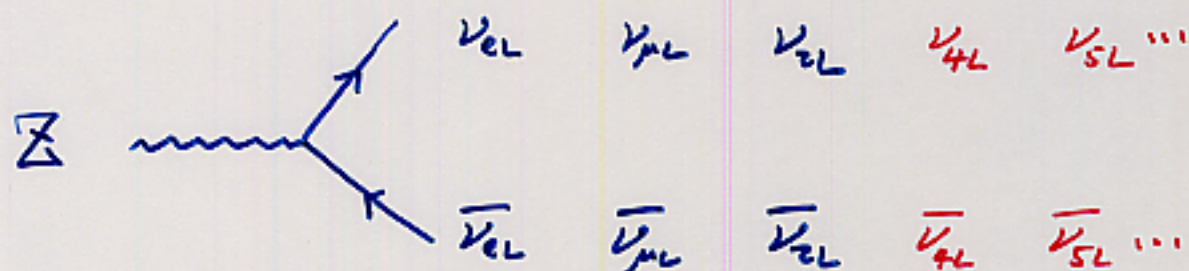
• 精密データを使っ易い形で保存

• 三世代のクォークとレプトン

第一世代: $\begin{pmatrix} u_L \\ d_L \end{pmatrix}$ u_R d_R $\begin{pmatrix} \nu_e \\ e_L \end{pmatrix}$ e_R

第二世代: $\begin{pmatrix} c_L \\ s_L \end{pmatrix}$ c_R s_R $\begin{pmatrix} \nu_\mu \\ \mu_L \end{pmatrix}$ μ_R

第三世代: $\begin{pmatrix} t_L \\ b_L \end{pmatrix}$ t_R b_R $\begin{pmatrix} \nu_\tau \\ \tau_L \end{pmatrix}$ τ_R



$\Gamma(Z \rightarrow \text{invisible mode}) \propto N_\nu$

$\sigma_h^0 = \frac{12\pi}{m_Z^2} \frac{\Gamma_e \Gamma_h}{\Gamma_Z^2}$

$R_\ell = \frac{\Gamma_h}{\Gamma_\ell} \quad (\ell = e, \mu, \tau)$

Γ_Z

m_Z

$\Rightarrow \Gamma_Z = \Gamma_h + \Gamma_\ell + \Gamma_{\text{inv.}}$

$N_\nu = 2.9841 \pm 0.0083$

* Degenerate 4'th generation with heavy Dirac ν_4 .

$\Rightarrow \Delta S = 0.21$

* Vector-like families

BB宇宙論と素粒子論

$$\frac{{}^4\text{He}}{\text{H}} \propto \frac{\#n}{\#p}$$

平衡 $\frac{\#n}{\#p} = e^{-\frac{m_n - m_p}{T}} \approx e^{-\frac{1.3\text{MeV}}{T}}$

↓
冷却 $G_F^2 T^5 \quad < \quad \sqrt{G_N (5.5 + \frac{7}{4} N_\nu)} T^2$
 $n \leftrightarrow p$ 宇宙膨張

$N_\nu \uparrow \quad T \uparrow \quad \frac{\#n}{\#p} \uparrow \quad \frac{{}^4\text{He}}{\text{H}} \uparrow$
観測値 ~ 0.25
↓
 $N_\nu = 2 \sim 4$



- * バリオン生成 (物質宇宙の起源)
- * ダークマター (銀河の全役)
- * ダークエネルギー・宇宙定数 (???)

弱相互作用のゲージ理論

自発的ゲージ対称性の破れ (ヒッグス機構)

→ $m_W, m_Z, m_\rho, m_\ell, \dots$

→ 短距離・高振動部は ゲージ対称性を保つ

↓

高エネルギー(未知の)物理に
対する依存性がゲージ対称性
により制限される。

↓

くり込み可能

↓

精密実験による検証。

v, g, g', \dots ($g'/g = \tan \theta_w$)

m_W

$\sim v g$



低振動部 ↔ 高振動部

は $E = \hbar \omega - \epsilon$

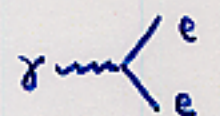
m_Z

$\sim v g / \cos \theta_w$

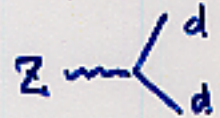


$G_F \sim \frac{g^2}{m_W^2}$

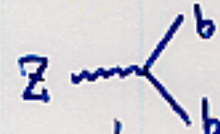
$\sim v^{-2}$



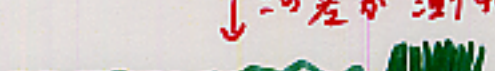
$\sim e = g \sin \theta_w$



$\sim \frac{g}{\cos \theta_w}$



$\sim \frac{g}{\cos \theta_w}$



↑ この差が測れる。

Precision Electroweak Data (Lepton-Photon 2001)

← QED

← QCD

	data	SM	pull	● 1998 data
LEP 1				
line-shape & FB asym.:				
m_Z (GeV)	91.1867 ± 0.0021	—	—	—
Γ_Z (GeV)	2.4952 ± 0.0023	2.4968	-0.7	2.4959
● σ_{had}^0 (nb)	41.540 ± 0.037	41.476	1.7 ●	41.473
R_e	20.804 ± 0.050	20.744	1.2	20.746
R_μ	20.785 ± 0.033	20.744	1.3	20.746
R_τ	20.764 ± 0.045	20.791	-0.6	20.793
$A_{\text{FB}}^{0,e}$	0.0145 ± 0.0025	0.0165	-0.8	0.0162
$A_{\text{FB}}^{0,\mu}$	0.0169 ± 0.0013	0.0165	0.3	0.0162
$A_{\text{FB}}^{0,\tau}$	0.0188 ± 0.0017	0.0165	1.4	0.0162
τ polarization:				
A_τ	0.1439 ± 0.0042	0.1482	-1.0	0.1471
A_e	0.1498 ± 0.0048	0.1482	0.3	0.1471
b and c quark results:				
R_b	0.21646 ± 0.00065	0.21597	0.8	0.21607
R_c	0.1719 ± 0.0031	0.1721	-0.1	0.1720
● $A_{\text{FB}}^{0,b}$	0.0990 ± 0.0017	0.1039	-2.9 ●	0.1031
● $A_{\text{FB}}^{0,c}$	0.0685 ± 0.0034	0.0743	-1.7 ●	0.0737
jet charge asymmetry:				
$\sin^2 \theta_{\text{eff}}^{\text{lept}}$	0.2321 ± 0.0012	0.2314	0.9	0.2315
SLC				
● A_{LR}^0	0.1513 ± 0.0021	0.1482	1.5 ●	0.1471
A_b	0.922 ± 0.020	0.935	-0.6	0.935
A_c	0.670 ± 0.026	0.668	0.1	0.668
Tevatron + LEP 2				
● m_W (GeV)	80.451 ± 0.033	80.397	1.6 ●	80.372
Parameters				
$1/\alpha(m_Z^2)$	128.936 ± 0.046	128.92	0.3	128.93
$\alpha_s(m_Z)$	0.118 ± 0.002	0.118	0.0	0.118
● m_t (GeV)	174.3 ± 5.1	176.3	-0.4	173.5
m_H (GeV)	—	119	—	138
χ_{Tot}^2 (d.o.f. = 22 - 4)			● 26.1	● 17.5

Systematic approach to study quantum corrections

K.H. Ann.Rev.Nucl.Part.Sci. 48, 463 (1998)

The Z-parameters

$Z \rightarrow f_\alpha \bar{f}_\alpha$ amplitudes

$$g_\alpha^f = \text{reference SM value} + \underbrace{(\dots) \Delta \bar{g}_Z^2 + (\dots) \Delta \bar{S}^2}_{\text{oblique corrections}} + \underbrace{\Delta g_\alpha^f}_{\text{vertex corrections}}$$

$$\begin{cases} \Delta \bar{g}_Z^2 = 0.00412 \Delta T_Z \\ \Delta \bar{S}^2 = 0.00360 \Delta S_Z - 0.00241 \Delta T_Z \end{cases}$$

$$\begin{cases} \Delta S_Z = \Delta S + \Delta R - 0.064 x_\alpha \\ \Delta T_Z = \Delta T + 1.49 \Delta R - \frac{\Delta \delta g}{\alpha} \end{cases}$$

$$x_\alpha = \frac{\frac{1}{\alpha(m_Z^2)} - 128.90}{0.09}$$

← corrections to μ -decay

Oblique parameters $\Delta S, \Delta T, \Delta R$

$$\uparrow \sim \frac{1}{\bar{g}_Z^2(m_Z^2)} - \frac{1}{\bar{g}_Z^2(0)}$$

m_t, m_H dependence

$$\Delta S = (\Delta S)_{SM} + S_{\text{new}}$$

$$\Delta T = (\Delta T)_{SM} + T_{\text{new}}$$

$$\Delta R = (\Delta R)_{SM} + R_{\text{new}}$$

} oblique corrections

$$\Delta g_\alpha^f = (\Delta g_\alpha^f)_{SM} + (\Delta g_\alpha^f)_{\text{new}} \quad \text{vertex corrections}$$

only $\Delta g_L^b \neq 0$
in SM

12 distinctive corrections in MSSM

$\nu_e, e_L, e_R, u_L, u_R, d_L, d_R$
 $\nu_\tau, \tau_L, \tau_R, b_L, b_R$

The W-mass $m_W = 80.402 - 0.288 \Delta S + 0.418 \Delta T + 0.337 \Delta R + 0.012 x_\alpha - 0.126 \frac{\Delta \delta g}{\alpha}$

oblique corrections

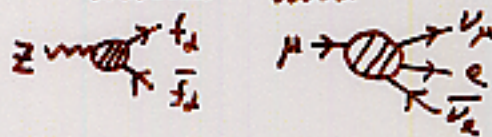
$Z \rightarrow f_\alpha \bar{f}_\alpha$ amplitudes

$$\begin{aligned}
 g_L^\nu &= 0.50214 + 0.453 \Delta \bar{J}_Z^2 && + \Delta g_L^\nu \\
 g_L^e &= -0.26941 - 0.244 \Delta \bar{J}_Z^2 + 1.001 \Delta \bar{S}^2 && + \Delta g_L^e \\
 g_R^e &= 0.23201 + 0.208 \Delta \bar{J}_Z^2 + 1.001 \Delta \bar{S}^2 && + \Delta g_R^e \\
 g_L^\mu &= 0.34654 + 0.314 \Delta \bar{J}_Z^2 - 0.668 \Delta \bar{S}^2 && + \Delta g_L^\mu \\
 g_R^\mu &= -0.15466 - 0.139 \Delta \bar{J}_Z^2 - 0.668 \Delta \bar{S}^2 && + \Delta g_R^\mu \\
 g_L^d &= -0.42451 - 0.383 \Delta \bar{J}_Z^2 + 0.334 \Delta \bar{S}^2 && + \Delta g_L^d \\
 g_R^d &= 0.07732 + 0.069 \Delta \bar{J}_Z^2 + 0.334 \Delta \bar{S}^2 && + \Delta g_R^d \\
 g_L^b &= -0.42109 - 0.383 \Delta \bar{J}_Z^2 + 0.334 \Delta \bar{S}^2 && + \Delta g_L^b
 \end{aligned}$$

SM : $(S, T, U, R, \Delta g_L^b)$ ~ functions $(m_t, m_H, \alpha, \alpha_s)$

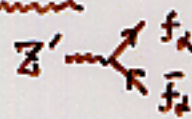
MSSM : $(S, T, U, R, 12 \Delta g_\alpha^f, \Delta \delta G)$ ~ functions (MSSM parameters)

UPDATE
G.C.Cho + KH
hep-ph/01



G.C.Cho + KH
NPB574, 623 (2000)
used CERN-EP/99-15
Data

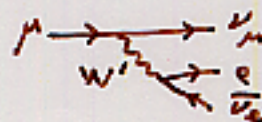
Z' models : $(+ T_{\text{new}}, \Delta g_\alpha^f)$
Z-Z' mixing



Triplet Higgs v.e.v.
: (T_{new})

G.C.Cho, Y.Umeda + KH
NPB531, 65 (1998)

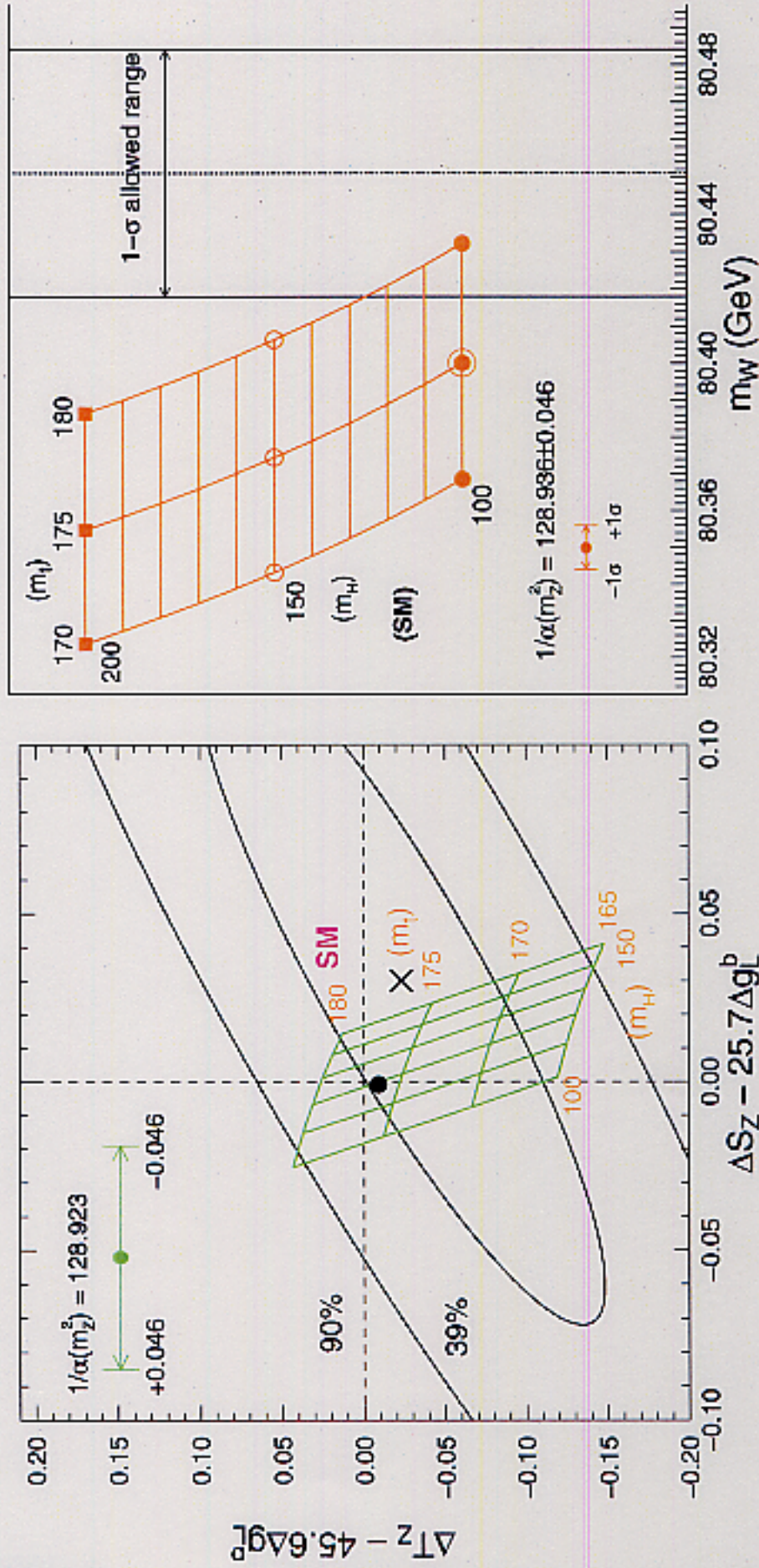
W' models : $(+ U_{\text{new}}, \Delta \delta G)$
W-W' mixing



KH + S. Matsumoto
PLB 424, 362 (1998)
E 8498, 456 (2000)

TC models : (S, T, U, \dots)

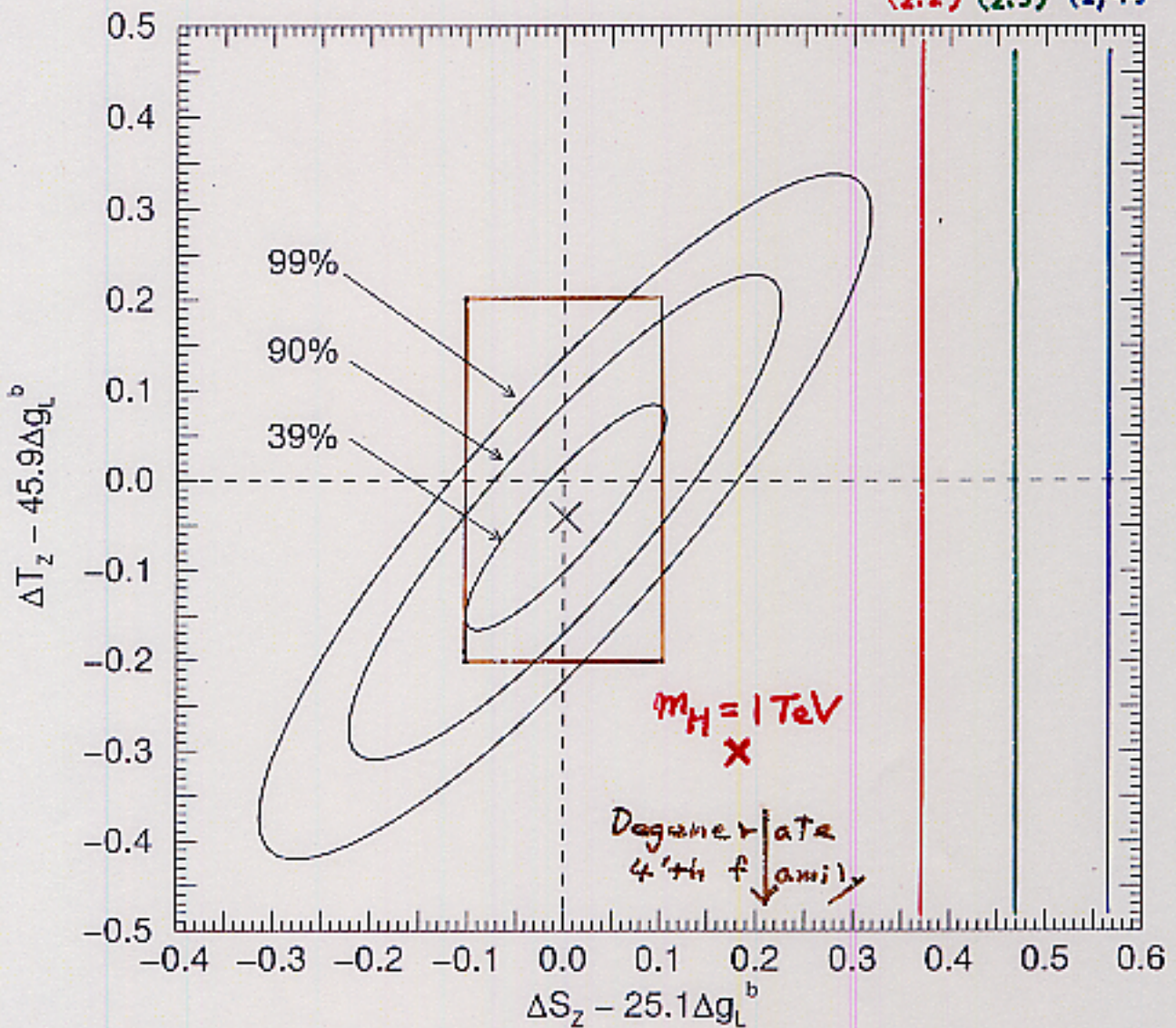
The SM prediction



QCD-like TC theory

(N_{TF}, N_{TC})

(2,2) (2,3) (2,4)



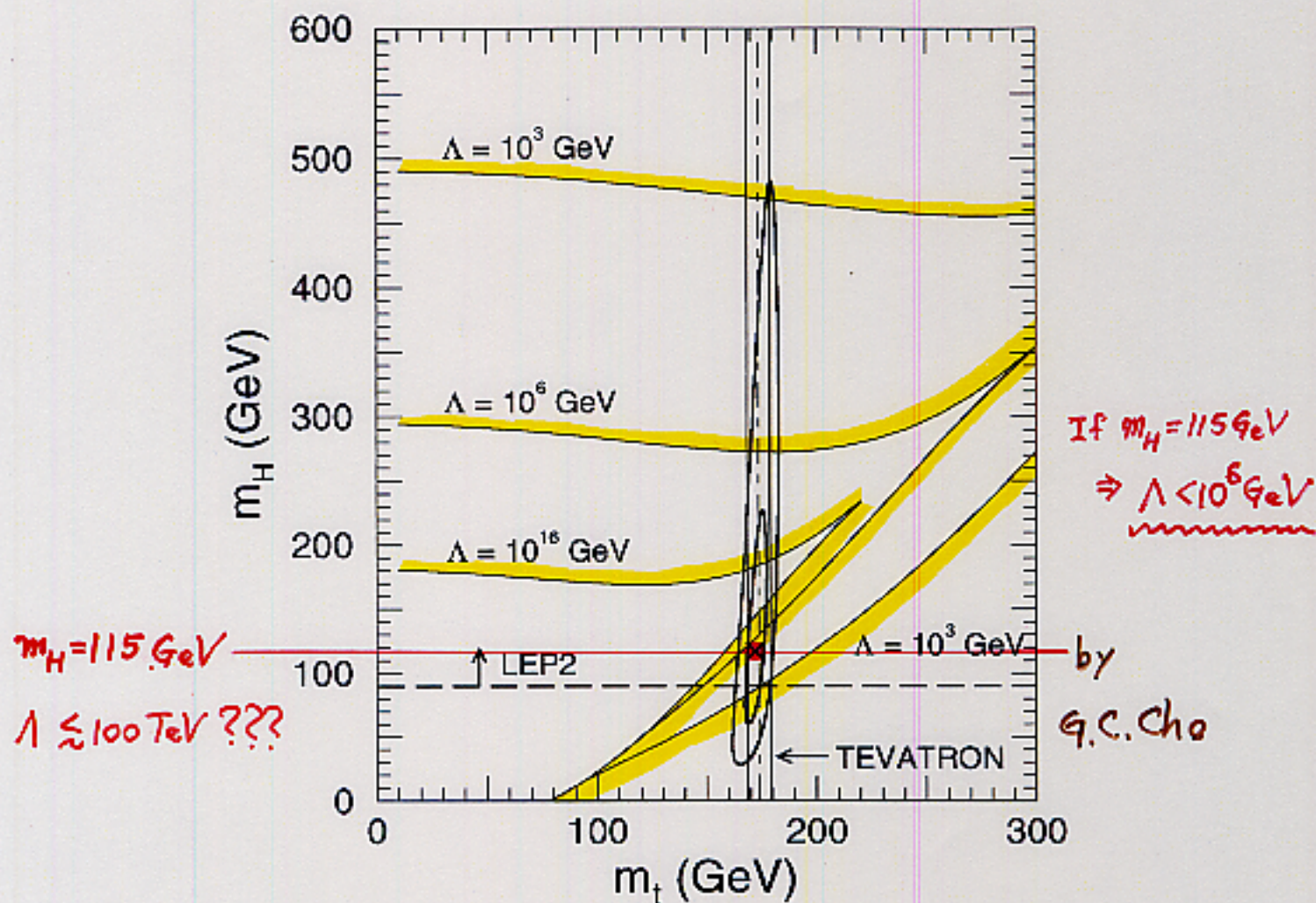
If $m_H = 1 \text{ TeV}$

If TC-like theory

If 4th generation

} \Rightarrow additional contribution
with fine-tuning
necessary.

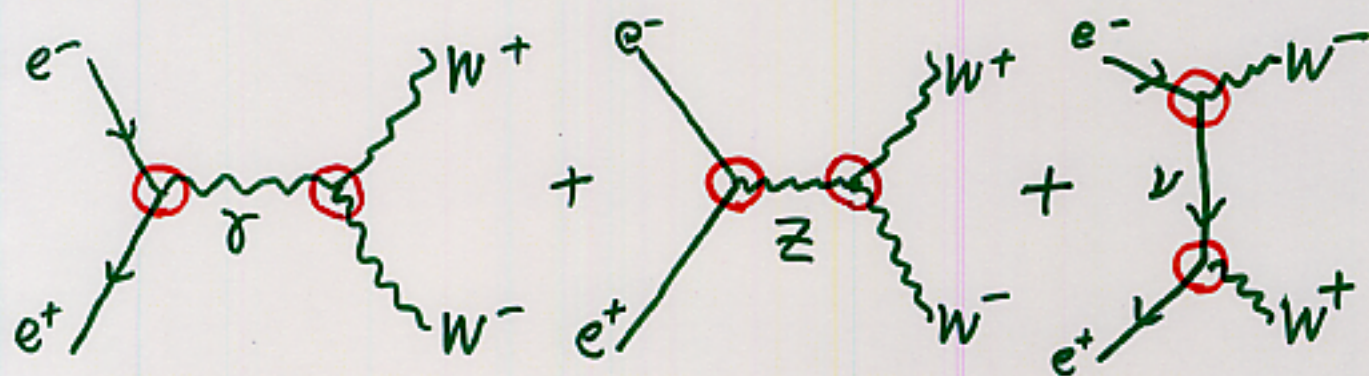
\Rightarrow A set of new particles.



90% (39%) CL allowed region from precision experiments at LEP, SLC, Tevatron (Sep., 1999). Also shown are the triviality (upper) bounds and the vacuum stability (lower) bounds for the Higgs boson mass.

$m_H \lesssim 200 \text{ GeV}$ favored by the electroweak precision measurements at LEP1/SLC, m_W at Tevatron/LEP2

WW γ /WWZ 結合の測定



ゲージ対称性 \Rightarrow WW γ /WWZ結合がユ= -1

\Downarrow

くり込み可能

\Downarrow

輻射補正によるテスト
(精密実験)

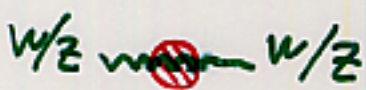

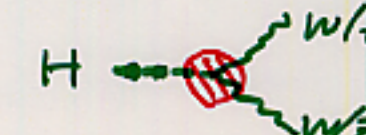
\Downarrow

直接測定 @ LEP2

\Downarrow

新しい物理探索のモデル

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \sum_i \frac{f_i}{\Lambda^2} \mathcal{O}_i^{(6)} + \mathcal{O}\left(\frac{1}{\Lambda^4}\right)$$

- 
精密電弱実験
- 
WWZ結合の精密測定
- 
HVV結合の精密測定

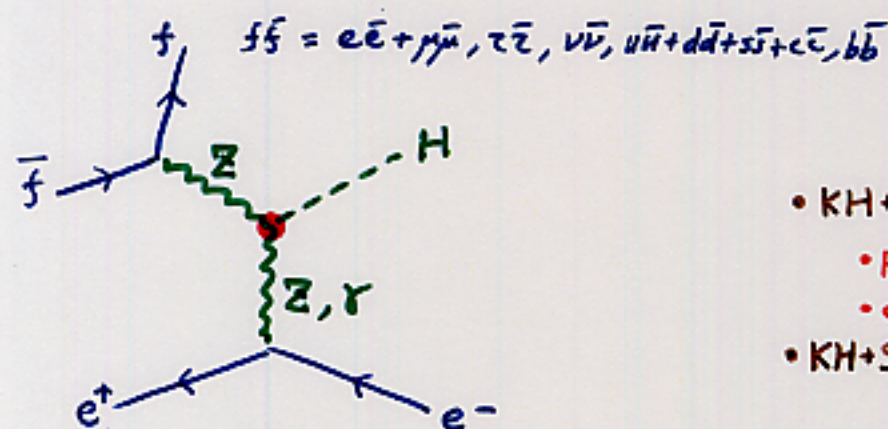
How accurately can we measure

the HVV couplings at a linear collider?

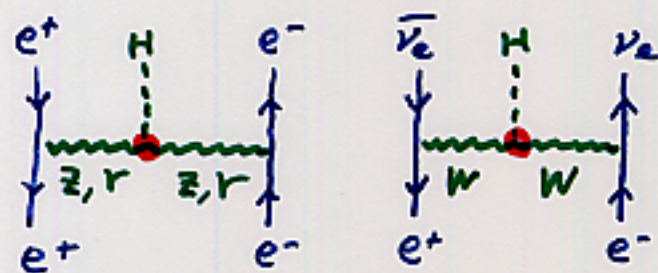
Kaoru Hagiwara (KEK)

2001.10.31

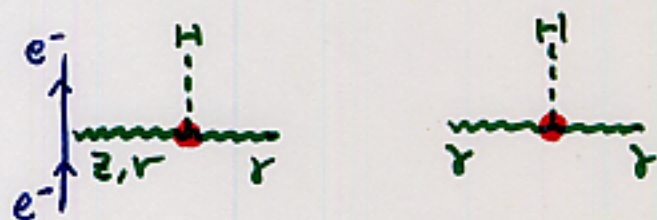
Tsinghua U., ACFA LCWS IV



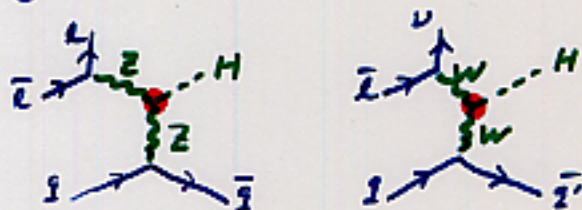
- KH + ML Stong, ZPC 62, 99 (1994)
 - parametrization of HVV couplings
 - observables
- KH + S Ishihara + JKamoshita + JAKniehl, EPJC 14, 457 (2000)
 - optimal observables \rightarrow error estimate
 - beam polarization, Z pol., b-tagging etc



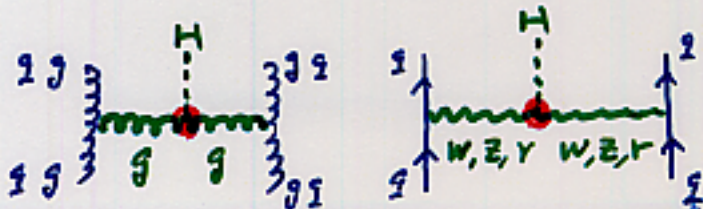
- KH + JKamoshita + Y Uehara, in preparation
 - tagged e^+, e^- information
 - beam polarization
- 2001.10



- Bremsstrahlung vs PLC



- Tevatron



- LHC

Table 6

Energy-dimension-six operators in the linear representation of the Higgs mechanism. The contribution of an operator to a particular vertex is denoted by an "X". In some cases an operator naively contributes to a vertex, yet that contribution does not lead to observable effects. In such cases the "X" is replaced by an "O"

$O_i^{(6)}$	WW	ZZ	AZ	AA	WWZ	WWA	WWWW	WWZZ	WWZA	WWAA	ZZZZ	HH	HW	HZ	HZ	HY	HY
$O_{DW} = \text{Tr}(D_\mu \hat{W}_{\nu\rho} D^\mu \hat{W}^{\nu\rho})$	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X
$O_{DB} = -\frac{g^2}{2} (\partial_\mu B_{\nu\rho}) (\partial^\mu B^{\nu\rho})$	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X
$O_{BW} = \phi^\dagger \hat{B}_{\mu\nu} \hat{W}^{\mu\nu} \phi$	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X
$O_{\phi,1} = (\partial_\mu \phi)^\dagger \phi (\partial^\mu \phi) $	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X
$O_{WWW} = \text{Tr}(\hat{W}_{\mu\nu} \hat{W}^{\nu\rho} \hat{W}^{\rho\mu})$	O	O	O	O	O	O	O	O	O	O	O	O	X	X	X	X	X
$O_{WW} = \phi^\dagger \hat{W}_{\mu\nu} \hat{W}^{\mu\nu} \phi$	O	O	O	O	O	O	O	O	O	O	O	O	X	X	X	X	X
$O_{BB} = \phi^\dagger \hat{B}_{\mu\nu} \hat{B}^{\mu\nu} \phi$	O	O	O	O	O	O	O	O	O	O	O	O	X	X	X	X	X
$O_W = (D_\mu \phi)^\dagger \hat{W}^{\mu\nu} (D_\nu \phi)$	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X
$O_B = (D_\mu \phi)^\dagger \hat{B}^{\mu\nu} (D_\nu \phi)$	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X
$O_{\phi,2} = \frac{1}{2} \partial_\mu (\phi^\dagger \phi) \partial^\mu (\phi^\dagger \phi)$	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X
$O_{\phi,3} = \frac{1}{3} (\phi^\dagger \phi)^3$	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X	X
$O_{\phi,4} = (\phi^\dagger \phi) (\partial_\mu \phi)^\dagger (D^\mu \phi) $	O	O	O	O	O	O	O	O	O	O	O	O	X	X	X	X	X

WWZ/8

大統一理論と超対称性のキート

何故 $Q_p + Q_e = 2Q_u + Q_d + Q_e = 0$?
 何故 $\alpha_s > \alpha_w > \alpha$?

} \Rightarrow 大統一理論
 \downarrow

大砂漠

\downarrow

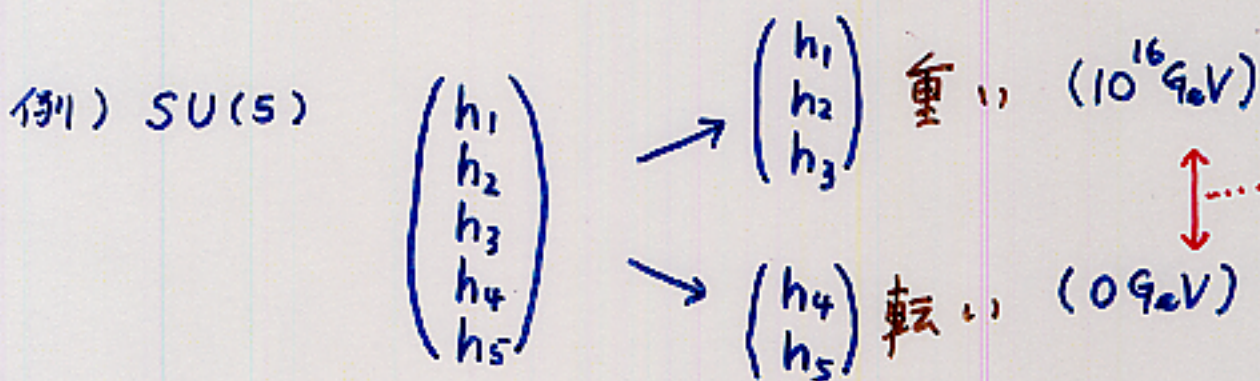
軽いスカラー場 (ヒッグス) の不安定性

\downarrow

超対称性

クォークとレプトン は同じものの部分

グルオンと W, Z と光 は同じものの部分



24 のゲージボソンの内、 $\begin{pmatrix} h_1 \\ h_2 \\ h_3 \end{pmatrix}$ に結合する部分

$\leftarrow g_1, g_3, B, \text{光子}$

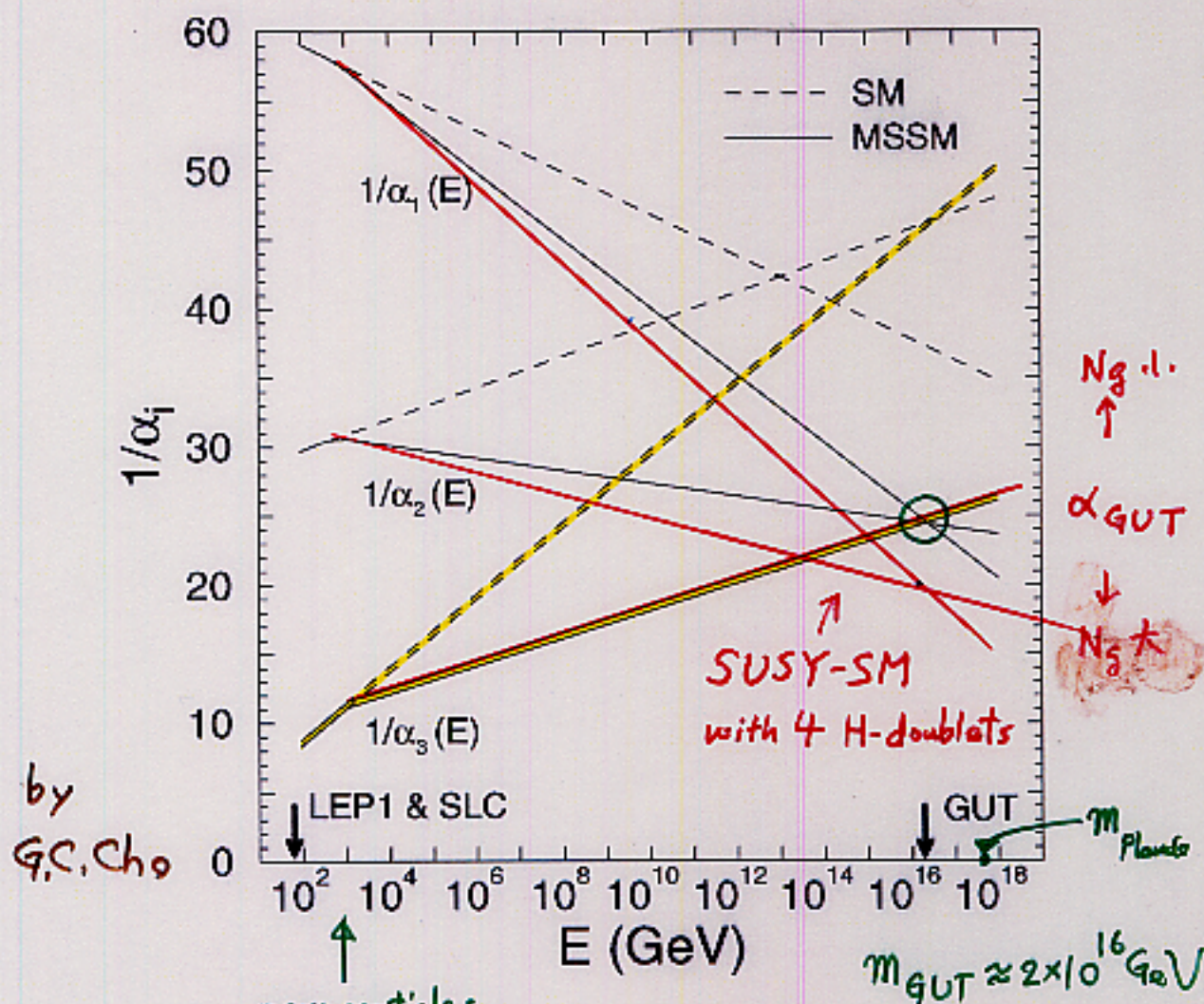
24 のゲージボソンの内、 $\begin{pmatrix} h_4 \\ h_5 \end{pmatrix}$ に結合する部分

$\leftarrow g_2, W, Z$

$\leftarrow \frac{g_1}{2}, \frac{g_2}{2}$

量子効果が種々異なる 10^{16} GeV

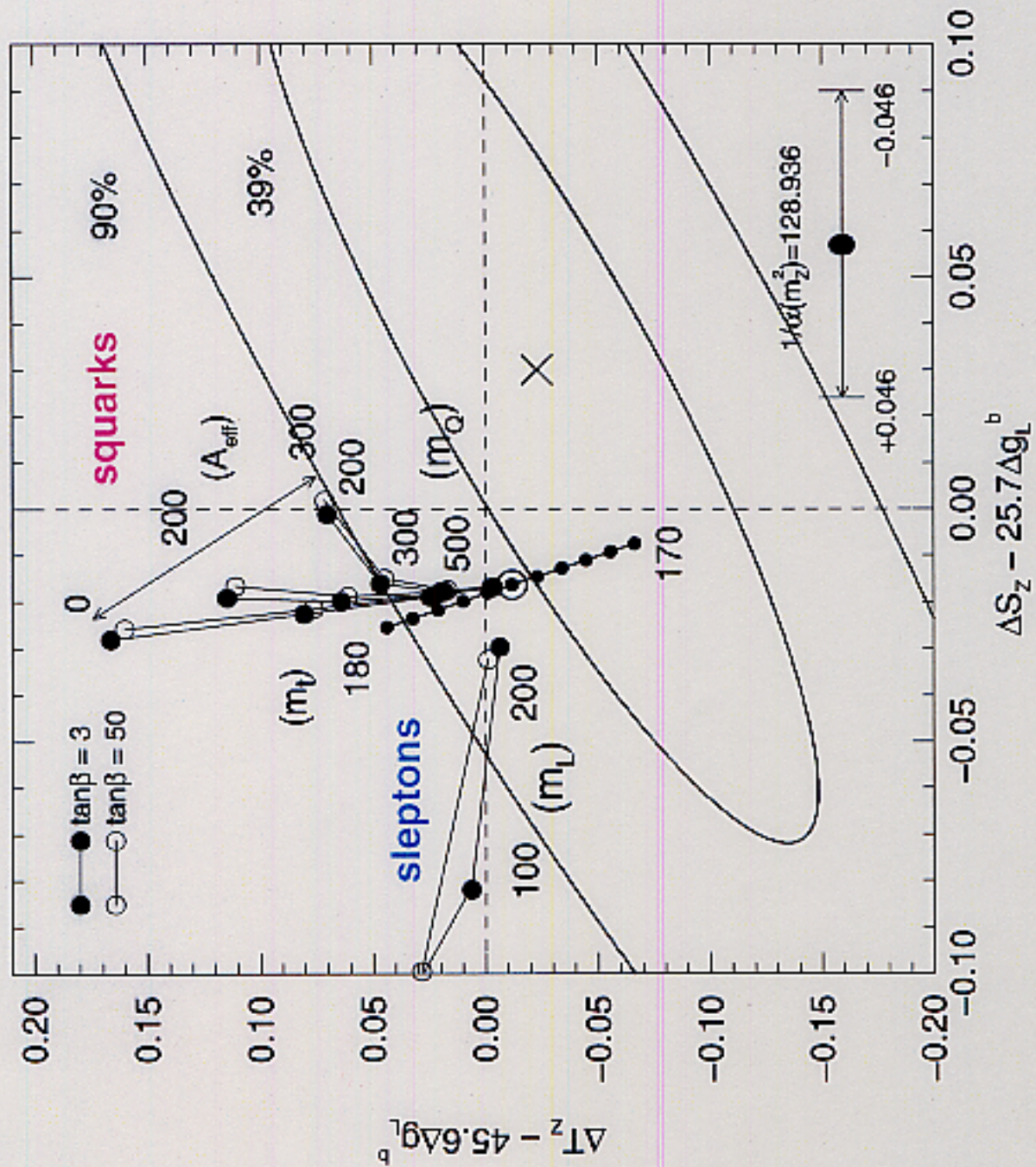
Running gauge couplings in the SM and in the MSSM (Minimal Supersymmetric SM)



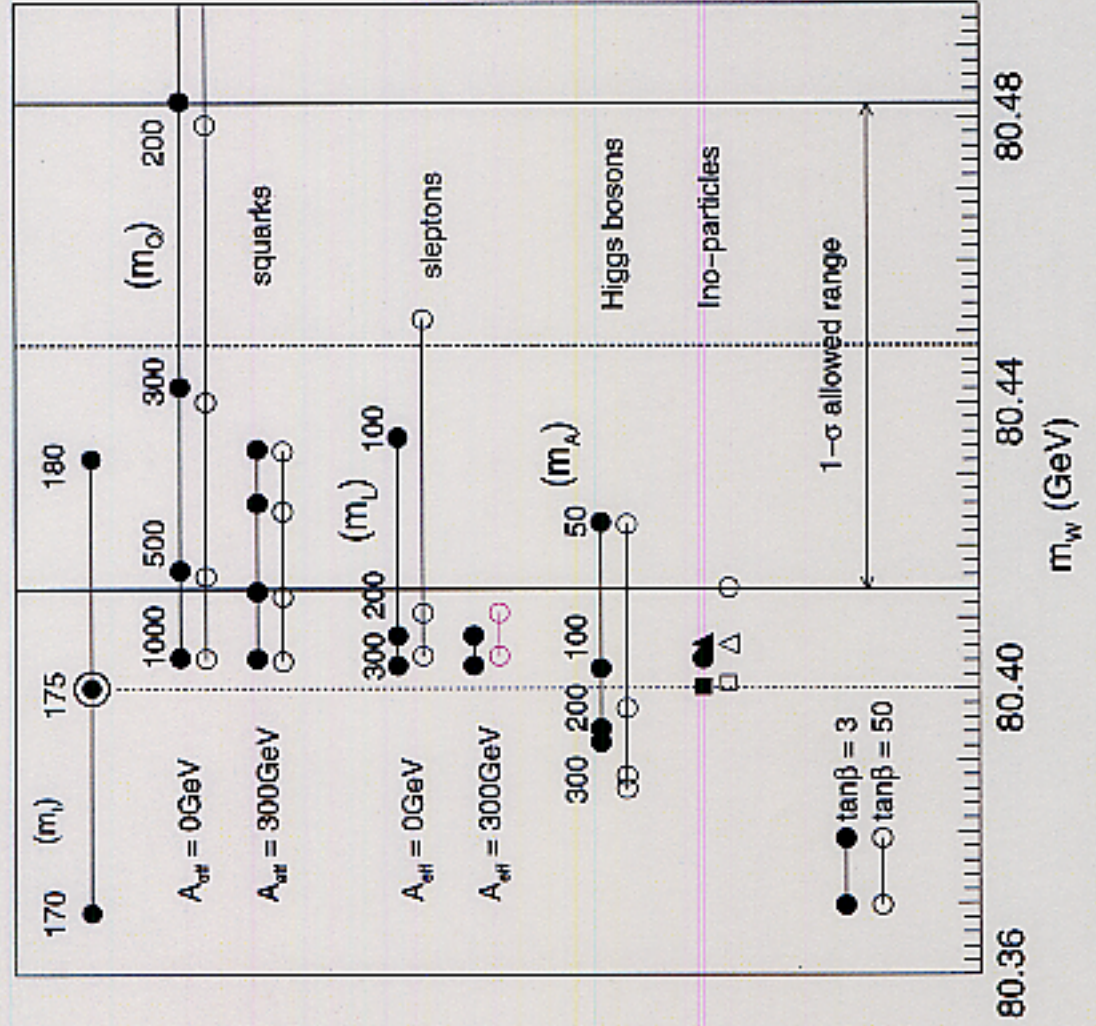
MSSM with 2 H-doublets
is favored! FCNC OK.

If some of them are sufficiently light,
their effects may be detected by precision
electroweak measurements of Z and W .

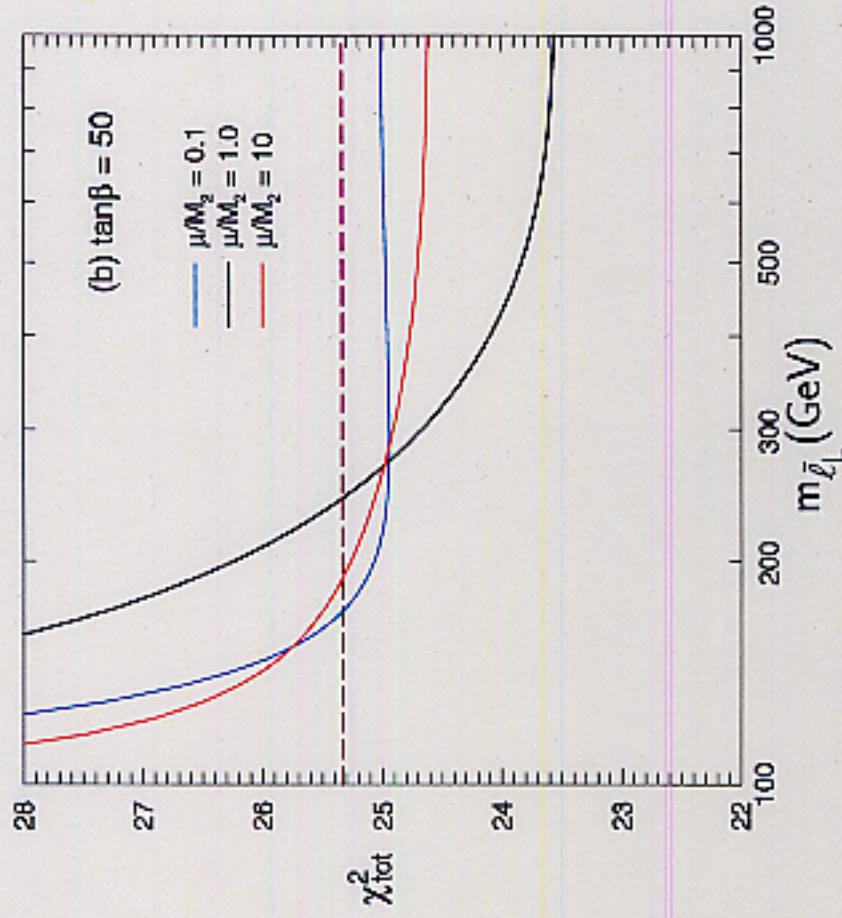
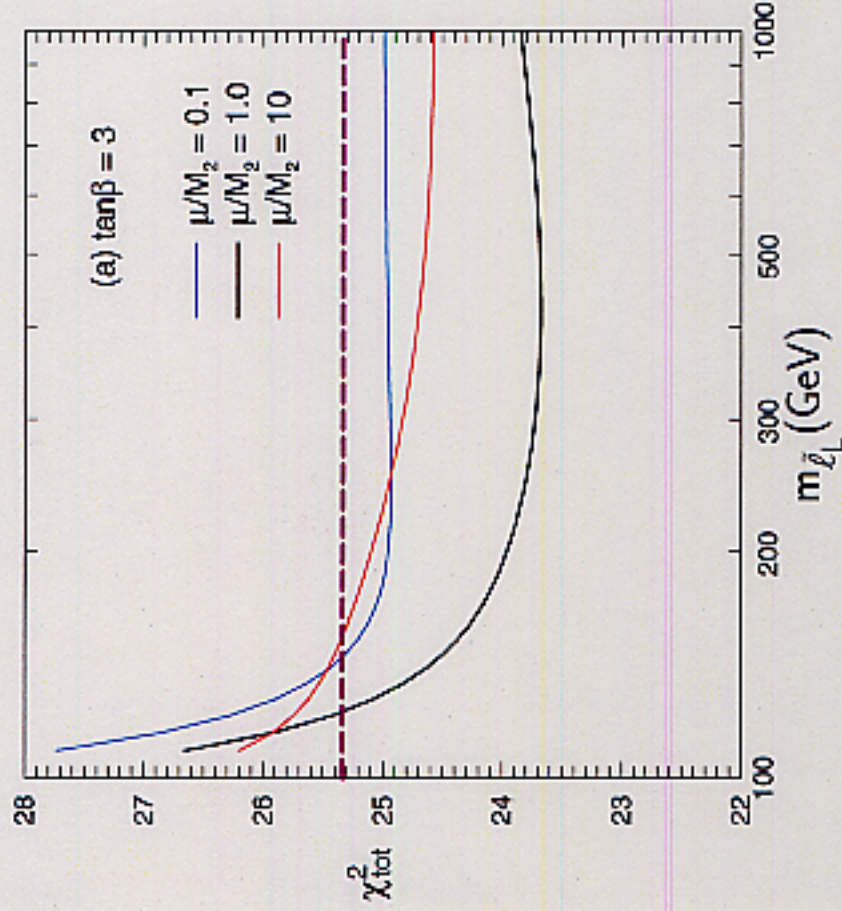
Sfermions



W-boson mass

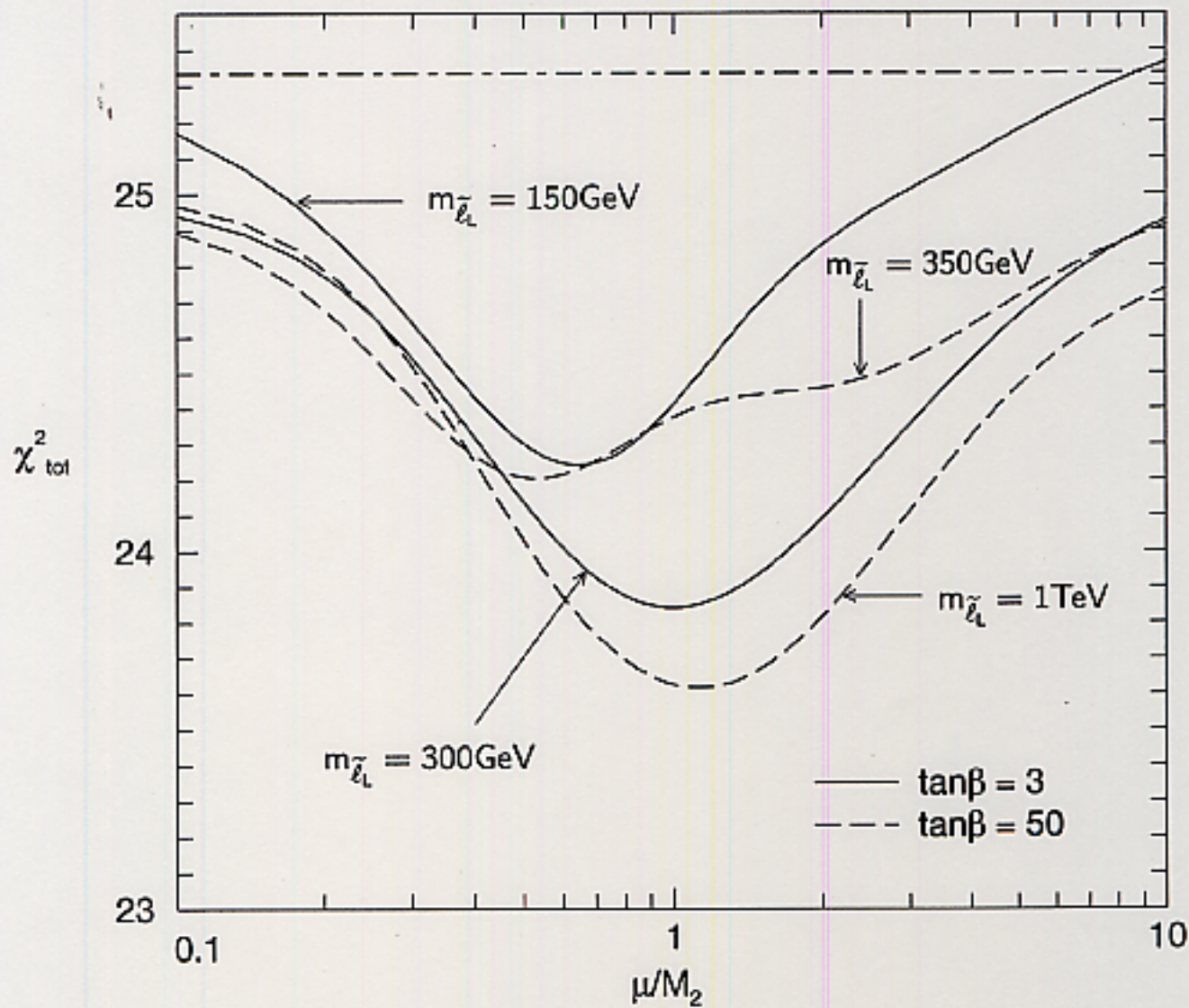


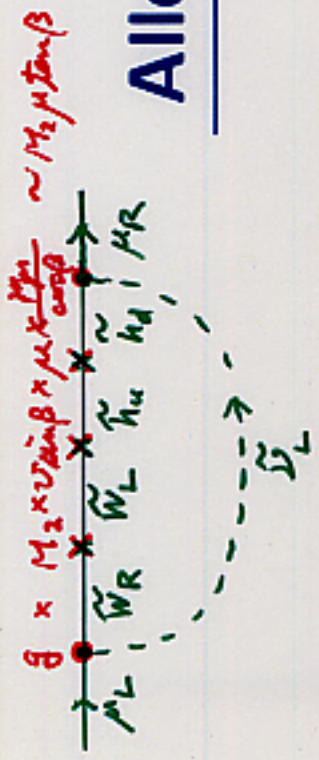
Light chargino & slepton



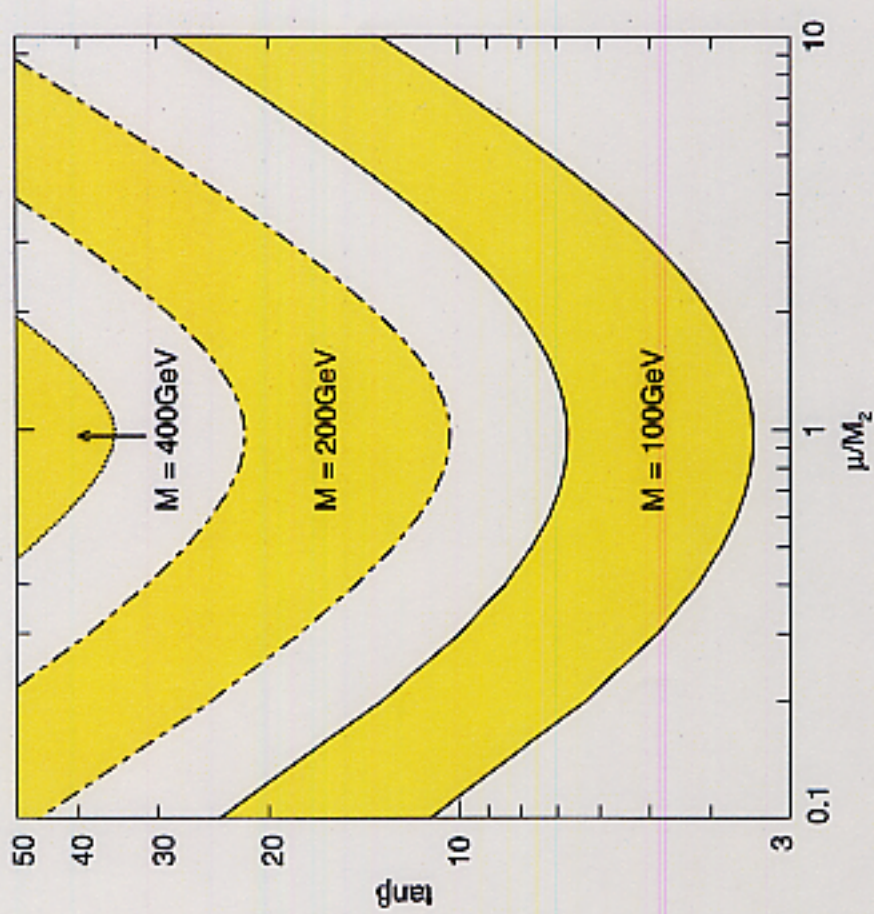
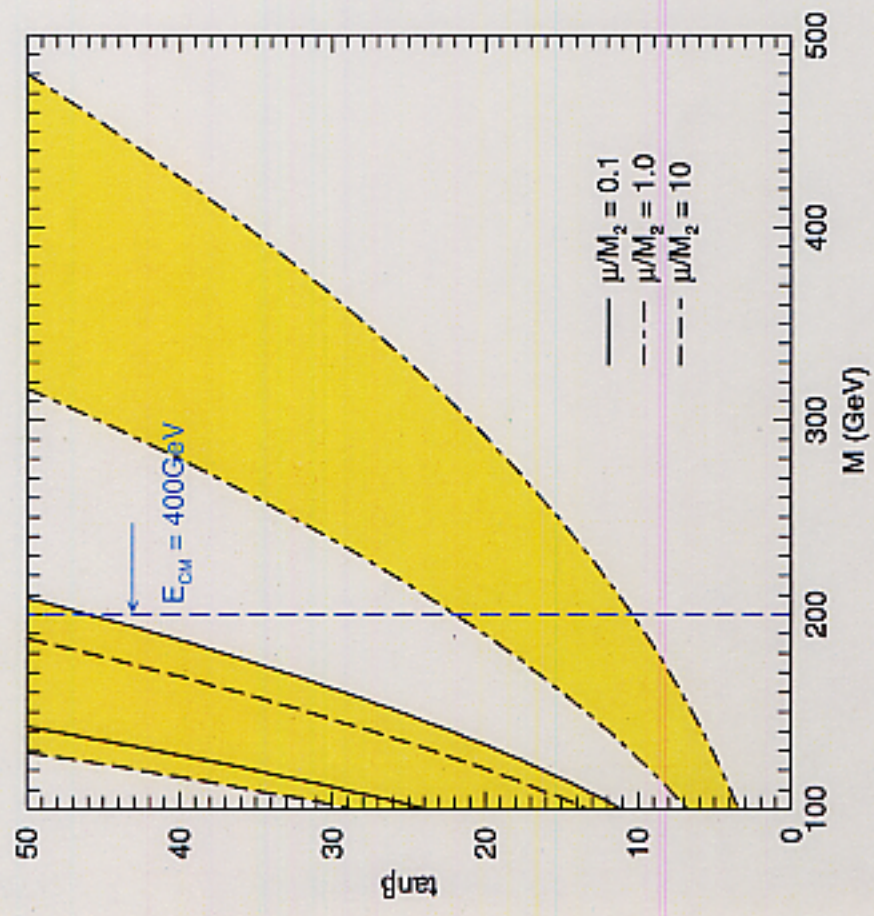
$\mu/M_2 \sim 1$ is favored from the EW precision data

$$\chi^2(\text{SUSY}) < \chi^2(\text{SM})$$





Allowed region [2] from μ g-2



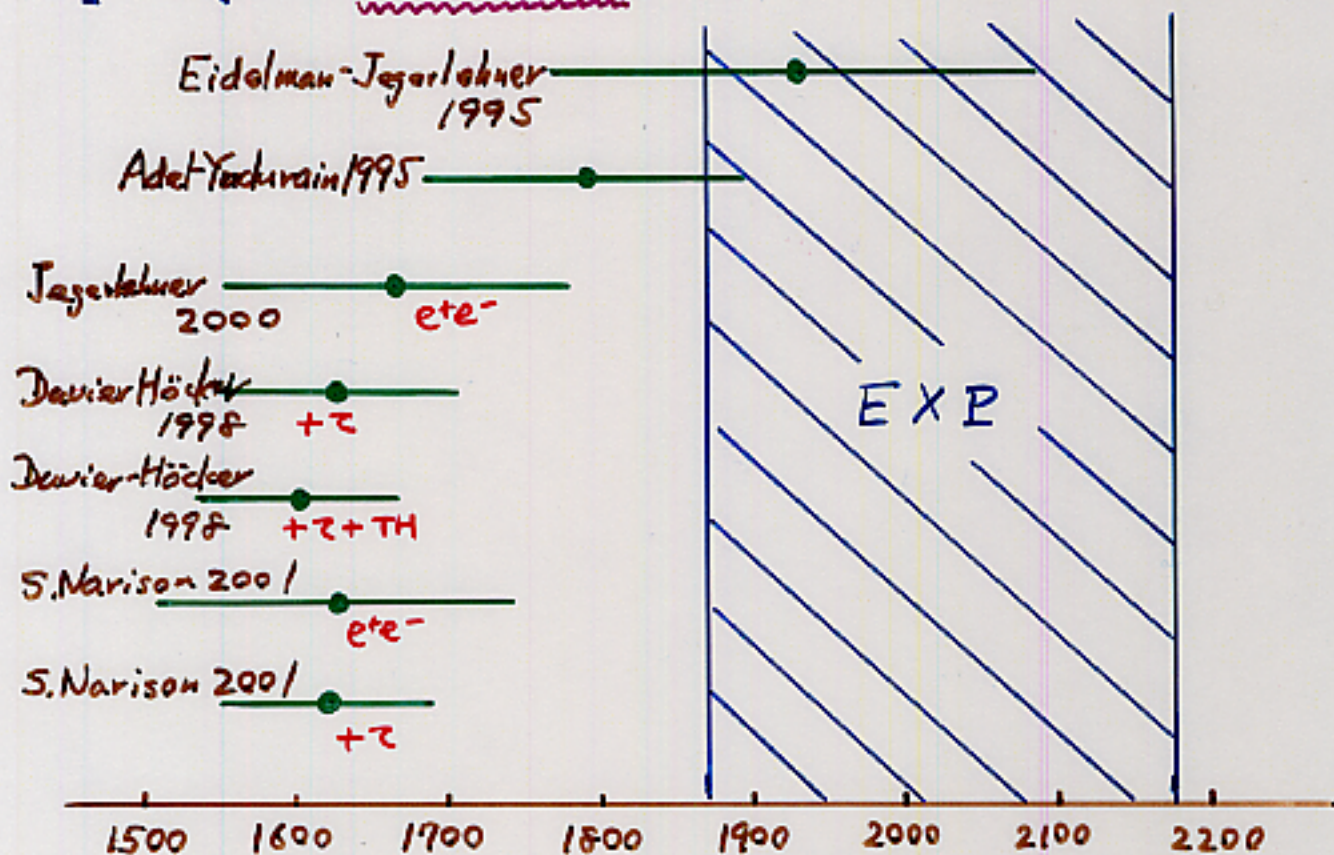
case for $m_{\tilde{X}_1^-} = m_{\tilde{\mu}_L} = m_{\tilde{\mu}_R}$

GCCo + KH, hep-ph/0105037
PLB

$(g-2)_\mu$

BNL: E821; H.N. Brown et al, hep-ex/0102017

$$\frac{g-2}{2} = (116\,592\,023 \pm 151) \times 10^{-11}$$



- Recent theoretical estimates are consistent.
- Use of τ -decays in $I=1$ channel + CVC reduces error.
error? \rightarrow Babar BELLE physics
- E821 error may be reduced to half soon...

Evidence or hint of New Physics?

$$\Delta a_\mu [\text{EXP-SM}] \approx 400 \pm 150 \times 10^{-11}$$

精密データの保存

LEP Electroweak Working Group

<http://lepewwg.web.cern.ch/LEPEWWG/>

LEPの4実験グループが共同作業により、
個々の実験データを総合し、模型依存性を
最小限に留める形で公表し、保存した。



- 外。研究者が独自の解析を行える。
- 新しい理論による解析ができる。
- 新現象との整合性が将来にわたって検証できる。

LEPの精密測定と新しい物理

2001. 11. 5

LEP ツンボツウム @ 東大

KEK 萩原 薫

- $N_\nu = 3$
 - \Rightarrow 3世代のクォーク・レプトン \Rightarrow フレーバーの物理
 - \Rightarrow 宇宙論と素粒子論 \Rightarrow ダークマター
ダークエネルギー
物質宇宙の起源
- 弱い相互作用のゲージ理論を検証
 - 輻射補正 $\Rightarrow m_t = 150 \sim 200 \text{ GeV} \Rightarrow$ 一致!
 - \Rightarrow W ボソンは2重項 \Rightarrow 素粒子模型
 - $\Rightarrow 113 \text{ GeV} \lesssim m_{H_{SM}} \lesssim 200 \text{ GeV} \Rightarrow$ 新しい物理を示唆
制限
 - \Rightarrow 「ゲージ対称性を自発的に破る機構の解明」
が素粒子実験の最大の課題となる。 \Rightarrow LHC, LC
- $WW\gamma$ / WWZ 結合の検証 \Rightarrow 「新しい物理」探索のモデル
- 大統一理論・超対称性のヒント
 - $\alpha_1(\mu) = \alpha_2(\mu) = \alpha_3(\mu)$ at $\mu \sim 10^{16} \text{ GeV}$
IF $m_{SUSY} \lesssim 1 \text{ TeV}$
 - $m_H \lesssim 200 \text{ GeV}$ は SUSY-GUT の予言と合う
 - $m_t \sim 175 \text{ GeV}$ は「輻射補正による対称性の破れ」の3行3
と合う
 - $m_{SUSY} < 1 \text{ TeV}$ のシグナル? $\Rightarrow (g-2)_\mu$?
- 精密データを使った易い形で保存
 - 今後何10年も「新しい物理のテスト」の場であり続けるだろう。